

## Optical Wired/Wireless Communication Using Soliton Optical Tweezers

S. E. Alavi<sup>1</sup>, I. S. Amiri<sup>2</sup>, S. M. Idrus<sup>1</sup>, J. Ali<sup>2</sup>

<sup>1</sup> Lightwave Communication Research Group, Faculty of Electrical Engineering, Universiti Teknologi Malaysia, 81310 UTM Skudai, Johor, Malaysia

<sup>2</sup> Institute of Advanced Photonics Science, Nanotechnology Research Alliance, Universiti Teknologi Malaysia (UTM), 81310 Johor Bahru, Malaysia

**Abstract:** The aim of this study is to generate nano optical tweezers to be connected to an optical quantum signal processing system in order to transmit quantum photon via an optical communication link. A system of microring resonator (MRR) known as Half-Panda is proposed to generate nano optical tweezers. Optical tweezers can be used to transport molecules in a communication link. The dark soliton propagates inside nonlinear MRR. Transportation of molecules or photons is implemented when the dark soliton is used as input pulse. The input Gaussian soliton is used to control the output signal at the through and drop ports of the system. Output nano optical tweezers can be connected to the quantum signal processing system consisting of a receiver and transmitter. The receiver will detect the signals of optical tweezers and transmit them via wired/wireless as quantum bits. The transmitter will generate the entangled photon pair which propagates via an optical communication link. Here the smallest nano optical tweezers signals with width at half maximum (FWHM) of 4.2 nm is obtained where the free spectrum range (FSR) of 50 nm is simulated.

[S. E. Alavi, I. S. Amiri, S. M. Idrus, J. Ali. **Optical Wired/Wireless Communication Using Soliton Optical Tweezers.** *Life Sci J* 2013;10(12s):179-187]. (ISSN:1097-8135). <http://www.lifesciencesite.com>. 33

**Keywords:** Half-Panda, Nano optical tweezers, quantum entangled photon

### 1.0 INTRODUCTION

Dark-Gaussian soliton controls within a semiconductor add/drop multiplexer has numerous applications in optical communication (Amiri *et al.* 2010a; I. S. Amiri and J. Ali). Nano optical tweezers technique has become a powerful tool for manipulation of micrometer-sized particles/photons in three spatial dimensions (Amiri and Ali 2012; Sadegh Amiri *et al.* 2013). It has the unique ability to trap and manipulate molecules/photons at mesoscopic scales with widespread applications in biology and physical sciences (I. S. Amiri and Ali 2013; Amiri and Ali 2013b). The output is achieved when the high optical field is set up as an optical tweezers (Afroozeh *et al.* 2010d; Imran *et al.* 2010; Nikoukar *et al.* 2013). For communication's application purposes, the optical tweezers can be used to generate entangled photon within the proposed network system (Amiri *et al.* 2013c; Nikoukar *et al.* 2012; Kouhnavard *et al.* 2010c; Amiri *et al.* 2012f).

The tweezers in the forms of valleys or potential wells are kept in the stable form within the add/drop filter (Amiri and Ali 2013c; Ali *et al.* 2010v; Amiri *et al.* 2010b). MRR's are type of Fabry-Perot resonators which can be readily integrated in array geometries to implement many useful functions (Amiri *et al.* 2011d; Shahidinejad *et al.* 2012; Nikoukar *et al.* 2010-2011). Its nonlinear phase response can be readily incorporated into an interferometer system to produce specific intensity output function (Amiri *et al.* 2012b; Afroozeh *et al.*

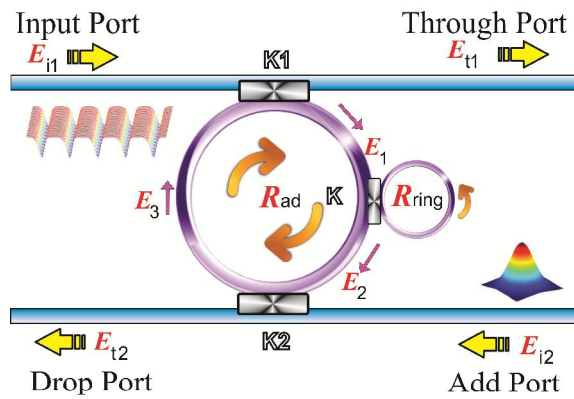
2010c; I. S. Amiri and J. Ali 2014c). Several emerging technologies, such as integrated all optical signal processing and all-optical quantum information processing (Kouhnavard *et al.* 2010a; Ali *et al.* 2010s; Ali *et al.* 2010y), require interactions between two distinct optical signals (Amiri *et al.* 2013e; Amiri *et al.* 2011a; I. S. Amiri *et al.* 2013b).

Internet security becomes an important function in the modern internet service (Amiri *et al.* 2012l; Amiri *et al.* 2011b). However, the security technique known as quantum cryptography has been widely used and investigated in many applications, using optical tweezers (Amiri *et al.* 2012g; Amiri *et al.* 2012n). Amiri *et al.* have proposed the new design of secured packet switching (Amiri *et al.* 2012d; I. S. Amiri *et al.* 2013d; I. S. Amiri *et al.* 2013e). This method uses nonlinear behaviors of light in MRR which can be used for high-capacity and security switching (I. S. Amiri and J. Ali 2014b; Afroozeh *et al.* 2011a). Recently quantum network shows promising usage for the perfect network security (Gifany *et al.* 2013; Amiri *et al.* 2010-2011). Amiri *et al.* have shown that the continuous wavelength can be generated by using a soliton pulse in a MRR (Amiri *et al.* 2013b; I. S. Amiri *et al.* ; Amiri and Ali 2013d). The secret key codes are generated via the entangled photon pair which is used to security purposes using the dark soliton pulse propagation (Amiri *et al.* 2011c; Amiri *et al.* 2012e; Shojaei and Amiri 2011a). In this study, a nano

molecular cryptography system based on optical soliton is developed.

## 2.0 THEORETICAL MODELING

The dark soliton pulse is introduced into the multiplexer half-Panda system shown in figure (1). This system consists of an add/drop filter system connected to a smaller ring resonator on the right side (I. S. Amiri *et al.* 2014; Amiri *et al.* 2012p). Dynamic behavior of the optical tweezers is appeared when the Gaussian soliton is input into the add port of the system. The ring resonator is connected to the add-drop interferometer system with radius ( $R_{ring}$ ) of 10  $\mu\text{m}$  and coupling coefficient ( $\kappa$ ) of 0.5. The effective area of the coupling section is  $A_{eff}=0.25 \mu\text{m}^2$ .



**Fig.1.** A schematic diagram of Half-Panda system

The input optical field ( $E_{i1}$ ) of the dark soliton and add optical field ( $E_{i2}$ ) of the Gaussian pulses are given by (Amiri and Ali 2014b; Amiri *et al.* 2012k; Afroozeh *et al.* 2012b)

$$E_{i1} = A \tanh\left[\frac{T}{T_0}\right] \exp\left[\left(\frac{z}{2L_D}\right) - i\omega_0 t\right], \quad (1)$$

$$E_{i2}(t) = A \exp\left[\left(\frac{z}{2L_D}\right) - i\omega_0 t\right] \quad (2)$$

In equations (1) and (2),  $A$  and  $z$  are the optical field amplitude and propagation distance, respectively (Afroozeh *et al.* 2012a; Ali *et al.* 2010h; Ali *et al.* 2010j).  $T$  is defined as soliton pulse propagation time in a frame moving at the group velocity (Amiri *et al.* 2012i; Ali *et al.* 2010e; Amiri *et al.* 2010c; Ali *et al.* 2010b),  $T = t - \beta_1 \times z$ , where  $\beta_1$  and  $\beta_2$  are the coefficients of the linear and second order terms of Taylor expansion of the propagation constant (I. S. Amiri *et al.* 2013a; Shojaei and Amiri 2011b).

$L_D = T_0^2 / |\beta_2|$  represents the dispersion length of the soliton pulse (Afroozeh *et al.* 2012c; Ali *et al.* 2010z; Ali *et al.* 2010g). The carrier frequency of the soliton

is  $\omega_0$  (Kouhnavard *et al.* 2010b; Mohamad *et al.* 2010b; Ali *et al.* 2010l; Ali *et al.* 2010x). When a soliton pulse keeps its temporal width invariance as it propagates, it is called a temporal soliton (A. Shahidinejad *et al.* 2014).  $T_0$  is known for the intensity of soliton peak as  $(|\beta_2 / \Gamma T_0^2|)$  (Afroozeh *et al.* 2011b; Ali *et al.* 2010c; Ali *et al.* 2010u). A balance should be achieved between the dispersion length ( $L_D$ ) and the nonlinear length ( $L_{NL} = (1/\gamma\phi_{NL})$ ) (Amiri *et al.* 2013f), where  $\gamma$  and  $\phi_{NL}$  are the coupling loss of the field amplitude and nonlinear phase shift (Amiri *et al.* 2012o; Ali *et al.* 2010p; Ali *et al.* 2010o; Afroozeh *et al.* 2010b). They are the length scale over which dispersive or nonlinear effects makes the beam becomes wider or narrower (I. S. Amiri and J. Ali 2014a; Ali *et al.* 2010a; Ali *et al.* 2010r; Ali *et al.* 2010n). It means that the  $L_D = L_{NL}$  should be satisfied (Amiri *et al.* 2013a; Yupapin *et al.* 2010). Within the nonlinear medium, the refractive index ( $n$ ) changes according to given following equation (3) (Amiri and Ali 2013a; Amiri *et al.* 2012a),

$$n = n_0 + n_2 I = n_0 + \left(\frac{n_2}{A_{eff}}\right) P, \quad (3)$$

$n_0$  and  $n_2$  are the linear and nonlinear refractive indexes, respectively (Amiri *et al.* 2012j).  $I$  and  $P$  represent the optical intensity and optical power, respectively (I. S. Amiri *et al.* 2013c; Ridha *et al.* 2010a; Ali *et al.* 2010q; Ali *et al.* 2010f). The effective mode core area of the device is given by  $A_{eff}$  (Amiri and Ali 2014a). In figure (1), the resonant output is formed, thus, the normalized output of the light field is the ratio between the output and input fields  $E_{out}(t)$  and  $E_{in}(t)$  (Ali Shahidinejad *et al.* 2014; Suwanpayak *et al.* 2010). The output and input signals in each roundtrip of the ring resonator on the right side can be calculated using equation (4) (Afroozeh *et al.* 2012d; I. S. Amiri and J. Ali 2014d).

$$\left|\frac{E_{out}(t)}{E_{in}(t)}\right|^2 = (1-\gamma) \left[ 1 - \frac{(1-(1-\gamma)x^2)\kappa}{(1-x\sqrt{1-\gamma}\sqrt{1-\kappa})^2 + 4x\sqrt{1-\gamma}\sqrt{1-\kappa}\sin^2\left(\frac{\phi}{2}\right)} \right] \quad (4)$$

Therefore ring resonator can be comparable to a Fabry-Perot cavity (Jalil *et al.* 2010). It has an input and output mirror with a field reflectivity,  $(1-\kappa)$ , and a fully reflecting mirror (Afroozeh *et al.* 2010a; Saktioto *et al.* 2010c). Here  $\kappa$  is the coupling coefficient, and  $x = \exp(-\alpha L/2)$  represents a roundtrip loss coefficient (Saktioto *et al.* 2010b; Ridha *et al.* 2010b),  $\Phi_0 = kLn_0$  and  $\Phi_{NL} = kLn_2|E_{in}|^2$  are the linear and nonlinear phase shifts (Ali *et al.* 2010d; Mohamad *et al.* 2010a; Ali *et al.* 2010w),  $k = 2\pi/\lambda$  is the wave propagation number in a vacuum (Teeka *et al.* 2011; Saktioto *et al.* 2010a; Ali *et al.* 2010k).  $L$

and  $\alpha$  are a waveguide length and linear absorption coefficient, respectively (Jalil *et al.* 2011; Amiri *et al.* 2012c). In this work, the iterative method is inserted to obtain the needed results using equation (4). Two complementary optical circuits of a ring-resonator add-drop filter can be expressed by the equations (5) and (6) (Amiri and Ali 2014c; Ali *et al.* 2010t).

$$\left| \frac{E_t}{E_{in}} \right|^2 = \frac{(1-\kappa_1) - 2\sqrt{1-\kappa_1} \cdot \sqrt{1-\kappa_2} e^{\frac{\alpha}{2}L_{ad}} \cos(k_n L_{ad}) + (1-\kappa_2)e^{-\alpha L_{ad}}}{1 + (1-\kappa_1)(1-\kappa_2)e^{-\alpha L_{ad}} - 2\sqrt{1-\kappa_1} \cdot \sqrt{1-\kappa_2} e^{\frac{\alpha}{2}L_{ad}} \cos(k_n L_{ad})} \quad (5)$$

and

$$\left| \frac{E_d}{E_{in}} \right|^2 = \frac{\kappa_1 \kappa_2 e^{\frac{\alpha}{2}L_{ad}}}{1 + (1-\kappa_1)(1-\kappa_2)e^{-\alpha L_{ad}} - 2\sqrt{1-\kappa_1} \cdot \sqrt{1-\kappa_2} e^{\frac{\alpha}{2}L_{ad}} \cos(k_n L_{ad})}, \quad (6)$$

where  $E_t$  and  $E_d$  represent the optical fields of the throughput and drop ports respectively.  $\beta = kn_{eff}$  is the propagation constant (S. E. Alavi *et al.* 2013),  $n_{eff}$  is the effective refractive index of the waveguide and the circumference of the ring is  $L_{ad} = 2\pi R_{ad} \cdot R_{ad}$  is the radius of the ring. The phase constant can be simplified as  $\Phi = \beta L$  (Amiri *et al.* 2012e; Ali *et al.* 2010i; Saktioto *et al.* 2010d). The chaotic noise cancellation can be managed by using the specific parameters of the add-drop device in which required signals can be retrieved by the specific users (I. S. Amiri *et al.* ; Ali *et al.* 2011; Ali *et al.* 2010; Ali *et al.* 2010m). The waveguide (ring resonator) loss is  $\alpha = 0.5 \text{ dBmm}^{-1}$ . The fractional coupler intensity loss is  $\gamma = 0.1$  (Tanaram *et al.* 2011; Amiri *et al.* 2012m). In the case of add-drop device, the nonlinear refractive index is neglected. The output fields,  $E_{t1}$  and  $E_{t2}$  at the throughput and drop parts of the Half-Panda are expressed by (Bahadoran *et al.* 2011; Amiri *et al.* 2012h)

$$E_{t1} = -x_1 x_2 y_2 \sqrt{\kappa_1} E_{t2} e^{\frac{\alpha L_{ad}}{2}} - j k_n \frac{L_{ad}}{2} + \left[ \frac{x_2 x_3 \kappa_1 \sqrt{\kappa_2} E_0 E_{t1} (e^{\frac{\alpha L_{ad}}{2}} - j k_n \frac{L_{ad}}{2})^2 + x_3 x_4 y_1 y_2 \sqrt{\kappa_1} \sqrt{\kappa_2} E_0 E_{t2} (e^{\frac{\alpha L_{ad}}{2}} - j k_n \frac{L_{ad}}{2})^3}{1 - x_1 x_2 y_1 y_2 E_0 (e^{\frac{\alpha L_{ad}}{2}} - j k_n \frac{L_{ad}}{2})^2} \right] \quad (7)$$

$$E_{t2} = x_2 y_2 E_{t2} + \left[ \frac{x_1 x_2 \kappa_1 \sqrt{\kappa_1} \sqrt{\kappa_2} E_0 E_{t1} e^{\frac{\alpha L_{ad}}{2} - j k_n \frac{L_{ad}}{2}} + x_1 x_3 y_1 y_2 \sqrt{\kappa_2} E_0 E_{t2} (e^{-\frac{\alpha L_{ad}}{2} - j k_n \frac{L_{ad}}{2}})^2}{1 - x_1 x_2 y_1 y_2 E_0 (e^{\frac{\alpha L_{ad}}{2}} - j k_n \frac{L_{ad}}{2})^2} \right] \quad (8)$$

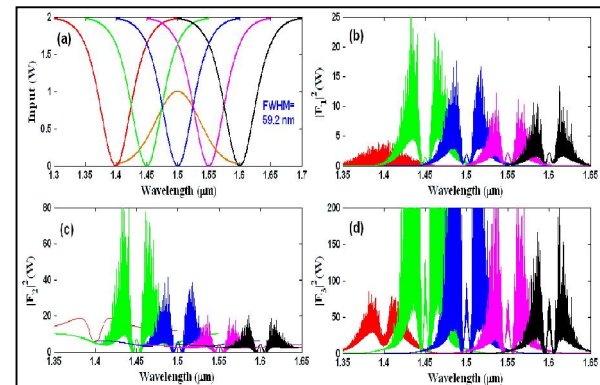
The electric field of the small ring on the right side of the Half-Panda system is given as (A. Nikoukar *et al.* 2014; Amiri *et al.* 2013d):

$$E_0 = E_1 \frac{\sqrt{(1-\gamma)(1-\kappa)} - (1-\gamma)e^{-\frac{\alpha}{2}L_{ring} - j k_n L_{ring}}}{1 - \sqrt{1-\gamma} \sqrt{1-\kappa} e^{-\frac{\alpha}{2}L_{ring} - j k_n L_{ring}}}, \quad (9)$$

where  $L_{ring} = 2\pi R_{ring}$  and  $R_{ring}$  is the radius of the ring and the  $x_1 = \sqrt{1-\gamma_1}$ ,  $x_2 = \sqrt{1-\gamma_2}$ ,  $x_3 = 1-\gamma_1$ ,  $x_4 = 1-\gamma_2$ ,  $y_1 = \sqrt{1-\kappa_1}$  and  $y_2 = \sqrt{1-\kappa_2}$ .

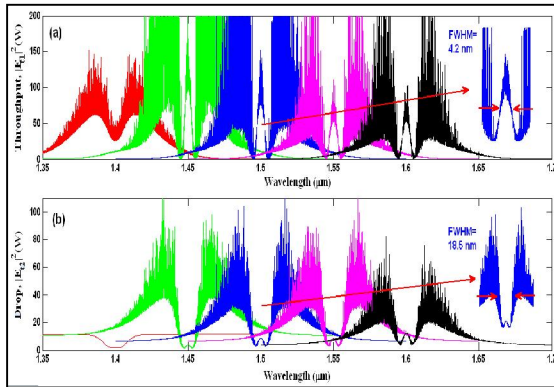
### 3.0 RESULT AND DISCUSSION

Input optical dark solitons and Gaussian laser beam with powers 2W and 1W respectively are inserted into the Half-Panda system. The add-drop optical filter has radius of  $R_{ad} = 15 \mu\text{m}$  where the coupling coefficients are  $\kappa_1 = 0.35$  and  $\kappa_2 = 0.25$ . The dark solitons are propagating inside the Half-Panda system with central wavelengths of  $\lambda_0 = 1.4 \mu\text{m}$ ,  $1.45 \mu\text{m}$ ,  $1.5 \mu\text{m}$ ,  $1.55 \mu\text{m}$ ,  $1.6 \mu\text{m}$ . In order to increase the capacity of the output optical tweezers, input dark soliton pulses with four different central wavelengths are introduced to the input port of the system. Figure 2(a), shows the optical inputs in the form of dark soliton and Gaussian laser beam. The nonlinear condition forms the interior signals as chaotic signals respect to 20,000 roundtrips of the input power. In order to make the system associate with the practical device (InGaAsP/InP), the selected parameters of the system are fixed to  $n_0 = 3.34$  and  $n_2 = 2.5 \times 10^{-17}$ . By adjusting the parameters such as the dark and Gaussian powers at the input and add ports and the coupling coefficients, the tweezers depth would be controlled and tuned as shown in figure 2(b-d). Amplification of the signals occurs within the nonlinear system which makes the signals suitable for long distance communication.



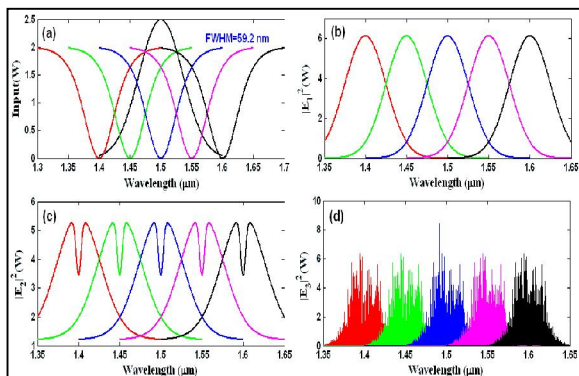
**Fig.2:** Optical tweezers generation within a Half-Panda system where (a): input of optical dark solitons and Gaussian laser beam, (b-d): amplified and tuned optical tweezers in the form of potential wells

Smallest tweezers width of 4.2 nm is generated at the through port shown in figure (3), where the drop port shows the output signal with FWHM of 18.5 nm. Transportation of the optical tweezers can be obtained via a network system using a photon generator and transmitter. Detection of the transported tweezers signal can be assembled using the single photon detection method. Thus, the tweezers transportation for long distance communication via molecular transporter is realistic.



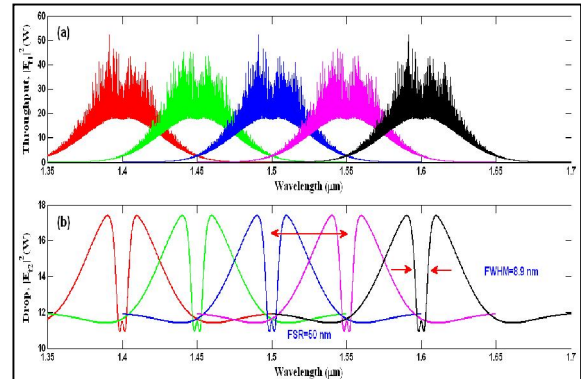
**Fig.3:** Through and drop port output signals of the Half-Panda system where (a): Through port output with FWHM=4.2 nm, (b): drop port output with FWHM=18.5 nm

Cancelling of the chaotic signals can be obtained within the add-drop ring resonator interferometer system using suitable parameters of the rings. The signals can be controlled and tuned by power's variation of the input Gaussian laser pulse. Figure (4) shows the generation of nano optical tweezers (clear signals). Here the input powers of the optical dark soliton pulses and Gaussian laser beam are 2W and 2.5W respectively.



**Fig.4:** Optical tweezers generation within a Half-Panda system where (a): input of dark solitons and Gaussian laser beam, (b-d): tuned and controlled optical tweezers

Filtered and clear optical tweezers are seen in figure (5) where the peaks have FWHM and FSR of 8.9 nm and 50 nm respectively. In the case of communication networks, generation of narrower signals is recommended. Therefore soliton signals can be used in optical communication where the capacity of the output signals can be improved by generation of peaks with smaller FWHM. The sensitivity of the ring systems can be improved significantly by generation of peaks with wider space or bigger FSR.



**Fig.5:** Through and drop port output signals of the Half-Panda system where (a): Through port chaotic output signals (b): drop port output with FWHM=8.9 nm and FSR=50 nm

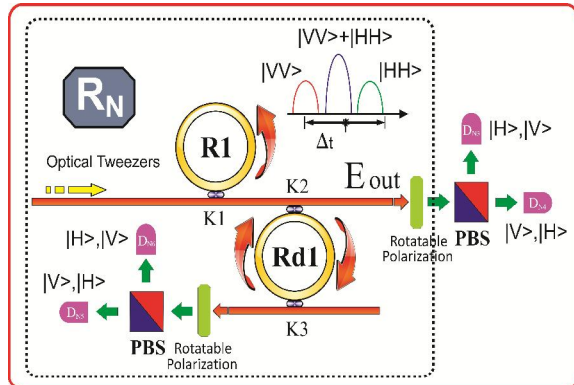
The proposed transmission unit is a quantum processing system that can be used to generate high capacity packet of quantum entanglement photons within the series of MRRs. In operation, the computing data can be modulated and input into the system via a receiver unit which is encoded to the quantum signal processing system. The receiver unit can be used to detect the quantum bits. It is obtained via the reference states recognized by using the cloning unit operated by an add-drop filter ( $R_{d1}$ ) shown in figure (6).

By using suitable dark-Gaussian soliton input power, tunable optical tweezers can be controlled. This provides the entangled photon as the dynamic optical tweezers probe. The required data can be retrieved via the through and drop ports of the add-drop filter in the router. High capacity data transmission can be applied by using more wavelength carriers. The advantage of this study is that ultra-short nano optical tweezers can be generated and transmitted via a network system thus improving the security and capacity.

The polarization states of light pulses are changed and converted during the circulation in the delay circuit, leading to the formation of the polarized entangled photon pairs. The entangled



photons of the nonlinear ring resonator are then separated into the signal and idler photon probability. The polarization angle adjustment device is applied to investigate the orientation and optical output intensity.



**Fig.6:** Schematic of entangled photon manipulation system. The quantum state is propagating to a rotatable polarizer and then is split by a beam splitter (PBS) flying to detector  $D_{N3}$ ,  $D_{N4}$ ,  $D_{N5}$  and  $D_{N6}$

Transporter states can be controlled and identified using the quantum processing system as shown in Figure 6. In operation, the encoded quantum secret codes computing data can be modulated and input into the system via a wireless router. Schematic of the wireless router is shown in Figure 7, in which quantum cryptography for internet security can be obtained. A wireless router system can be used to transfer generated entangled photons via a wireless access point, and network communication system shown in Figure 7.

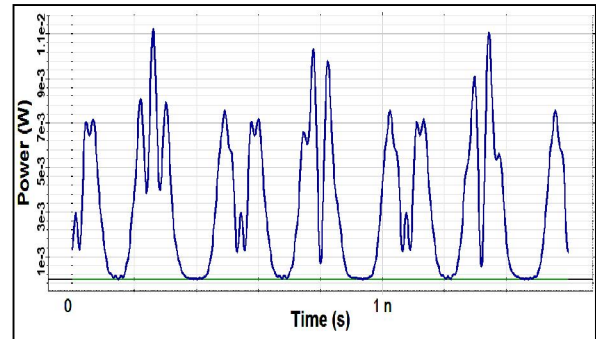


**Fig.7:** System of entangled photons transmission using a wireless access point system

A wireless access system transmits data to different users via wireless connection. This method also works in reverse, when the router system used to receive information from the Internet, translating it into an analog signal and sending it to the computer's wireless adapter. Figure 5 shows the detected and filtered optical tweezers using an optical receiver.

The advanced transmitter topologies are desirable for application in both wired and wireless communication inasmuch as they are able to provide power-efficient amplification of signals with large

peak-to average power ratios (PAPRs) without compromising system linearity.



**Fig.8:** Detected and filtered optical tweezers via a 25 m optical wireless link using an optical receiver

#### ACKNOWLEDGEMENTS

The authors acknowledge the administration of Universiti Teknologi Malaysia (UTM) and Research Management Centre (RMC) of UTM for the financial support through Geran Universiti Penyelidikan (GUP) funding vote number 04H35.

#### 4.0 CONCLUSION

Novel system of nano optical tweezers generation is presented for cryptography and secured optical communication. Nano optical tweezers can be generated by the dark soliton propagation in a Half-Panda system. The control process of the dark-Gaussian collision can be implemented by using suitable parameters of the ring system such as the input power, coupling coefficient, ring radius, coupler loss and effective core area. The Quantum signal processing unit is connected to the optical tweezers which can be used to generate quantum entanglement photons thus providing secured and high capacity information. This secured coded information can be easily transmitted via a communication network system. Here Double security can be performed when ultra-short of soliton pulses are used to generate entangle photon transmitting over long distance communication.

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