

Intrinsic formulation for elastic line deformed on a surface by an external field in the pseudo-Galilean space

$$G_1^3$$

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Abstract: In this paper, we derive intrinsic formulation for elastic line deformed on a surface by an external field in the pseudo-Galilean space G_1^3 [Nevin Gürbüz. **Intrinsic formulation for elastic line deformed on a surface by an external field in the pseudo-Galilean space** G_1^3 . *Life Sci J* 2013;10(4):1348-1352]. (ISSN:1097-8135). <http://www.lifesciencesite.com>. 178

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1.Introduction

Manning studied intrinsic formulation for elastic line deformed external field on a surface by external field E^3 (Manning, 1988). Intrinsic equations for a elastic line in Lorentz-Minkowski space was researched (Gürbüz and Görgülü, 2000), (Gürbüz, 2000). In this paper we derive intrinsic formulation for elastic line deformed external field on a surface by external field in pseudo-Galilean space.

In this section we give preliminaries on pseudo-Galilean space G_1^3 . The definitions relation to G_1^3 was taken (Divjak, 2008).

The pseudo-Galilean 3- space G_1^3 is the three dimensional real affine space with the absolute figure $\{w,f,I\}$, where w is a fixed plane, f a line in w and I a hyperbolic involution of the points of f . The pseudo-Galilean space length of the vector $x(x,y,z)$ is defined by

$$\begin{cases} x, & x \neq 0 \\ \sqrt{|y^2 - z^2|}, & x = 0 \end{cases}$$

A curve parametrized by the parameter of arc length $s=x$ is given in the coordinat form by $\beta(x)=(x,y(x),z(x))$. The curvature $\kappa(x)$ and $\tau(x)$ of an curve are given by (Divjak, 2008).

$$\kappa(x) = \sqrt{|y'^2(x) - z'^2(x)|}$$

$$\tau(x) = \frac{1}{\kappa^2(x)} \det(r'(x), r''(x), r'''(x))$$

The associated moving trihedron is given by

$$t = r'(x) = (1, y'(x), z'(x))$$

$$n = \frac{1}{\kappa(x)} (0, y''(x), z''(x)),$$

$$b = \frac{1}{\kappa(x)} (0, \varepsilon z''(x), \varepsilon y''(x))$$

where $\varepsilon = 1$ or $\varepsilon = -1$ and it is called a Frenet trihedron associated to the curve. If t is timelike, n is a spacelike vector, b is spacelike,. Frenet-Serret formulas are given as following:

$$t'(x) = \kappa(x)n(x)$$

$$n'(x) = \tau(x)b(x)$$

$$b'(x) = \tau(x)n(x).$$

For regular curve in G_1^3 , κ is defined as following

$$\kappa = \frac{\|\Psi' \times_{PG} \Psi''\|}{\|\Psi'\|^3}$$

where \times_{PG} denotes pseudo-Galilean cross product. If e_1 is unit spacelike vector, e_2 is unit spacelike vector, e_3 is a unit timelike vector, $a \times_{PG} b$ is given as following:

$$a \times_{PG} b = \begin{vmatrix} 0 & e_2 & -e_3 \\ a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \end{vmatrix}$$

where $a = (a_1, a_2, a_3), b = (b_1, b_2, b_3)$. If e_1 is unit spacelike vector, e_2 is unit timelike vector, e_3

is a unit spacelike vector , $a \times_{PG} b$ is given as following:

$$a \times_{PG} b = \begin{vmatrix} 0 & -e_2 & -e_3 \\ a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \end{vmatrix}$$

Theorem 1.1. Let F be the timelike surface in G_1^3 and β denote an arc on F. The analogue of the Frenet-Serret formulas in pseudo-Galilean 3-space G_1^3 is

$$\begin{bmatrix} T' \\ Q' \\ N' \end{bmatrix} = \begin{bmatrix} 0 & \kappa_g & \kappa_n \\ 0 & 0 & -\tau_g \\ 0 & \tau_g & 0 \end{bmatrix} \begin{bmatrix} T \\ Q \\ N \end{bmatrix} \tag{1.1}$$

where κ_g is the geodesic curvature , τ_g is the geodesic torsion, κ_n is the normal curvature.

$$\langle T, T \rangle = -1, \quad \langle N, N \rangle = 1, \quad \langle Q, Q \rangle = 1$$

Theorem 1.2. Let F be the spacelike surface in G_1^3 and β denote an spacelike arc on F. The analogue of the Frenet-Serret formulas in pseudo-Galilean 3-space G_1^3 is

$$\begin{bmatrix} T' \\ Q' \\ N' \end{bmatrix} = \begin{bmatrix} 0 & \kappa_g & -\kappa_n \\ 0 & 0 & -\tau_g \\ 0 & -\tau_g & 0 \end{bmatrix} \begin{bmatrix} T \\ Q \\ N \end{bmatrix} \tag{1.2}$$

where κ_g is the geodesic curvature , τ_g is the geodesic torsion, κ_n is the normal curvature. Also,

$$\langle T, T \rangle = 1, \quad \langle Q, Q \rangle = 1, \quad \langle N, N \rangle = -1$$

2. Intrinsic Method

In this section, we study intrinsic formulation for elastic line deformed on surface by an external field in pseudo-Galilean space G_1^3 .

The arc β is called elastic line if it is extremal for the variational problem of (2.1) within the family of all arcs of length l on non-null surface F having the same initial point and initial direction as β in the pseudo-Galilean space G_1^3 .

If elastic line is exposed to a static force field, it has a trajectory that minimizes the sum of its elastic

energy and its energy of interaction with the field in G_1^3 . The problem is to minimize the energy E ,

$$E = \int_0^l \left(\frac{1}{2} b \kappa^2 - \theta \phi \right) ds \tag{2.1}$$

$$E(t) = \frac{1}{2} b I_1(t) - \theta I_2(t)$$

among elastic lines with trajectories $\phi(u(s), v(s))$ of fixed length l and arc length , $0 \leq s \leq l$, contained pseudo-Galilean surface $\phi(u, v)$ in pseudo-Galilean space G_1^3 .

$-\theta$ is constant measuring the strength of the external field, $\phi(u, v)$ gives its shape and κ denotes elastic bending energy in the pseudo-Galilean 3-space .

The equilibrium trajectory are the extrema of the sum of stress and potential energies in G_1^3 . The path of the elastic line have to satisfy a differential equation, which is derived by variational methods on the pseudo-Galilean 3-space.

Assume β lies in a coordinat patch $\phi(u, v)$ of F . Thus β is given as $\beta(s) = \phi(u(s), v(s))$. Also, $T(s) = \beta'(s)$,

$$Q(s) = p(s)\phi_u + q(s)\phi_v \tag{2.2}$$

for suitable scalar functions p(s) and q(s). Define $\Psi(\sigma; t) = \phi(u(\sigma) + t\eta(\sigma), v(\sigma) + t\xi(\sigma))$

$$l = \int_0^{\lambda(t)} \sqrt{\left\langle \frac{\partial \Psi}{\partial \sigma}, \frac{\partial \Psi}{\partial \sigma} \right\rangle_{PG}} d\sigma$$

for $0 \leq \sigma \leq l$.

Case I. Intrinsic formulation for elastic line deformed on a timelike surface by an external field in the pseudo-Galilean space G_1^3 .

i. If T is timelike , Q and N are spacelike ,

$$\frac{\partial \Psi}{\partial \sigma} \Big|_{t=0} = T, \quad 0 \leq \sigma \leq l \tag{2.3}$$

$$\frac{\partial^2 \Psi}{\partial \sigma^2} \Big|_{t=0} = T' = \kappa_g Q + \kappa_n N \tag{2.4}$$

$$\frac{\partial \Psi}{\partial t} \Big|_{t=0} = T' = \mu Q \tag{2.5}$$

With second differentiation Equation (2.5) , we obtain

$$\left. \frac{\partial^2 \Psi}{\partial t \partial \sigma} \right|_{t=0} = T' = \mu' Q - \mu \tau_g N \tag{2.6}$$

Third differentiation Equation (2.5) gives

$$\left. \frac{\partial^3 \Psi}{\partial t \partial \sigma^2} \right|_{t=0} = (\mu'' - \mu \tau_g^2) Q - (2\mu' \tau_g + \mu \tau_g') N \tag{2.7}$$

$$\left. \frac{\partial \lambda}{\partial t} \right|_{t=0} = 0$$

Lemma 2.1. In pseudo-Galilean 3-space,

Proof.

Let $I_1(t) = \int \kappa^2 ds$ denote the total square curvature of the arc $0 \leq \sigma \leq \lambda(t)$. For $t \neq 0$, the total square curvature is

$$I_1(t) = \int_0^{\lambda(t)} \left\| \left\langle \frac{\partial \Psi}{\partial \sigma} \times_{PG} \frac{\partial^2 \Psi}{\partial \sigma^2}, \frac{\partial \Psi}{\partial \sigma} \times_{PG} \frac{\partial^2 \Psi}{\partial \sigma^2} \right\rangle \left\| \left\langle \frac{\partial \Psi}{\partial \sigma}, \frac{\partial \Psi}{\partial \sigma} \right\rangle_{PG} \right\|^{-5/2} d\sigma$$

Therefore

$$I_1'(0) = \int_0^{\lambda(0)} \frac{\left\langle \left. \frac{\partial^3 \Psi}{\partial t \partial \sigma} \right|_{t=0}, \left. \frac{\partial^2 \Psi}{\partial \sigma^2} \right|_{t=0} \right\rangle_{PG} \left\| \left\langle \left. \frac{\partial^2 \Psi}{\partial \sigma^2} \right|_{t=0}, \left. \frac{\partial^2 \Psi}{\partial \sigma^2} \right|_{t=0} \right\rangle_{PG} \right\|^{-5/2}}{\left\| \left\langle \left. \frac{\partial \Psi}{\partial \sigma} \right|_{t=0}, \left. \frac{\partial \Psi}{\partial \sigma} \right|_{t=0} \right\rangle_{PG} \right\|^{3/2} \left\langle \left. \frac{\partial^2 \Psi}{\partial \sigma^2} \right|_{t=0}, \left. \frac{\partial^2 \Psi}{\partial \sigma^2} \right|_{t=0} \right\rangle_{PG}} d\sigma \tag{2.9}$$

From (2.9), we obtain

$$I_1'(0) = \int_0^l (\mu'' \kappa_g - \mu \kappa_g \tau_g^2 - 2\mu' \kappa_n \tau_g - \mu \kappa_n \tau_g') ds \tag{2.10}$$

Using integration by parts

$$\int_0^l \mu' \kappa_n \tau_g ds = 2\mu(l) \kappa_n(l) \tau_g(l) - 2 \int_0^l \mu \kappa_n' \tau_g ds - 2 \int_0^l \mu \kappa_n \tau_g' ds \tag{2.11}$$

$$\int_0^l \mu'' \kappa_g ds = \mu'(l) \kappa_g(l) - \mu(l) \kappa_g'(l) + \int_0^l \mu \kappa_g'' ds \tag{2.12}$$

Using Equations (2.5) ,(2.6), (2.11),(2.12) ,we obtain

$$I_1'(0) = \int_0^l \mu (\kappa_g'' + 2\kappa_n' \tau_g + \kappa_n \tau_g' - \kappa_g \tau_g^2) ds + \mu'(l) \kappa_g(l) - \mu(l) \kappa_g'(l) - 2\mu(l) \kappa_n(l) \tau_g(l) \tag{2.13}$$

$$I_2(t) = \int_0^l \varphi ds \tag{2.14}$$

Differentiating of Equation (2.14) at t=0,

$$\frac{d\lambda}{dt} \Big|_{t=0} \int_0^{\lambda(t)} \left\| \left\langle \left. \frac{\partial \Psi}{\partial \sigma} \right|_{t=0}, \left. \frac{\partial \Psi}{\partial \sigma} \right|_{t=0} \right\rangle_{PG} \right\| + \int_0^l \frac{\left\langle \left. \frac{\partial \Psi}{\partial \sigma} \right|_{t=0}, \left. \frac{\partial \Psi}{\partial \sigma} \right|_{t=0} \right\rangle_{PG} \left\langle \left. \frac{\partial \Psi}{\partial \sigma} \right|_{t=0}, \left. \frac{\partial \Psi}{\partial \sigma} \right|_{t=0} \right\rangle_{PG}^{-1}}{\left\| \left\langle \left. \frac{\partial \Psi}{\partial \sigma} \right|_{t=0}, \left. \frac{\partial \Psi}{\partial \sigma} \right|_{t=0} \right\rangle_{PG} \right\|^{-1/2}} = 0$$

From (2.3) and (2.6), we obtain

$$\left. \frac{\partial \lambda}{\partial t} \right|_{t=0} = 0 \tag{2.8}$$

$$I'_2(0) = \int_0^l \left(\frac{\partial \varphi}{\partial t} \right) \Big|_{t=0} = \int_0^l \mu \left[p \left(\frac{\partial \varphi}{\partial u} \right) + q \left(\frac{\partial \varphi}{\partial v} \right) \right] ds \tag{2.15}$$

From (2.13) and (2.15), for all choices of the function $\mu(s)$, $E'(0) = 0$, the given timelike arc β must satisfy two boundary conditions and differential equation in pseudo-Galilean 3-space

$$\begin{aligned} \text{(BC1)} \quad \kappa'_g(l) &= 0 \\ \text{(BC2)} \quad \kappa'_g(l) &= -2\kappa_n(l)\tau_g(l) \end{aligned} \tag{2.16}$$

$$\text{(DE)} \quad b[\kappa''_g + 2\kappa'_n\tau_g + \kappa_n\tau'_g - \kappa_g\tau_g^2] - \theta \left[p \left(\frac{\partial \varphi}{\partial u} \right) + q \left(\frac{\partial \varphi}{\partial v} \right) \right] = 0.$$

Case II. Intrinsic formulation for elastic line deformed on a spacelike surface by an external field in the pseudo-Galilean space G_1^3 .

N is imelikee, T and Q are spacelike:

For $\left| \kappa_g^2 - \kappa_n^2 \right| = \kappa_g^2 - \kappa_n^2$, we have

$$\begin{aligned} I'_1(0) &= \int_0^l \mu(\kappa''_g + 2\kappa'_n\tau_g + \kappa_n\tau'_g + \kappa_g\tau_g^2) ds + \\ &\quad \mu'(l)\kappa_g(l) - \mu(l)\kappa'_g(l) + 2\mu(l)\kappa_n(l)\tau_g(l) \end{aligned} \tag{2.17}$$

For all choices of the function $\mu(s)$, $E'(0) = 0$, the given spacelike β arc must satisfy two boundary conditions and differential equation in pseudo-Galilean 3-space.

$$\begin{aligned} \text{(BC1)} \quad \kappa'_g(l) &= 0 \\ \text{(BC2)} \quad \kappa'_g(l) &= -2\kappa_n(l)\tau_g(l) \end{aligned} \tag{2.18}$$

$$\text{(DE)} \quad b[\kappa''_g + 2\kappa'_n\tau_g + \kappa_n\tau'_g + \kappa_g\tau_g^2] - \theta \left[p \left(\frac{\partial \varphi}{\partial u} \right) + q \left(\frac{\partial \varphi}{\partial v} \right) \right] = 0$$

Given spacelike β arc must satisfy two boundary conditions and differential equation in pseudo-Galilean 3-space.

$$\begin{aligned} \text{(BC1)} \quad \kappa'_g(l) &= 0 \\ \text{(BC2)} \quad \kappa'_g(l) &= 2\kappa_n(l)\tau_g(l) \end{aligned} \tag{2.19}$$

$$\text{(DE)} \quad b[\kappa''_g + 2\kappa'_n\tau_g + \kappa_n\tau'_g + \kappa_g\tau_g^2] + \theta \left[p \left(\frac{\partial \varphi}{\partial u} \right) + q \left(\frac{\partial \varphi}{\partial v} \right) \right] = 0$$

3. Results

Theorem 3.1. On the timelike surface in in pseudo-Galilean space for the case $\Theta = 0$, an timelike geodesic arc is elastic line if and only if it satisfies

$$\kappa_n^2 \tau_g = 0 \tag{2.20}$$

Since $\kappa_g = 0$, from the third equation of (2.16),

$$-2\kappa'_n\tau_g - \kappa_n\tau'_g = 0 \tag{2.21}$$

From (2.21), first integral is obtained

$\kappa_n^2 \tau_g = \text{constant}$. The constant must vanish, from the second equation of (2.16).

Theorem 3.2. An timelike geodesic arc on the timelike surface in pseudo-Galilean space for the case $\Theta \neq 0$, is elastic line if and only if it satisfies

$$b\kappa_n^2 \tau_g - \Theta \left[p \left(\frac{\partial \varphi}{\partial u} \right) + q \left(\frac{\partial \varphi}{\partial v} \right) \right] = 0 \tag{2.22}$$

Proof. From (2.16), we get (2.22).

Theorem 3.3. . An spacelike geodesic arc on the spacelike surface in pseudo-Galilean space for the case $\Theta \neq 0$, is elastic line if and only if it satisfies

$$b\kappa_n^2 \tau_g - \Theta \left[p \left(\frac{\partial \varphi}{\partial u} \right) + q \left(\frac{\partial \varphi}{\partial v} \right) \right] = 0 \tag{2.23}$$

Proof. From (2.18), we have (2.23).

Example 3.1. An timelike arc on timelike plane for $\Theta = 0$ in G_1^3 , is elastic line if and only if it lies on a geodesic.

Proof. On timelike plane, $\tau_g = (k_2 - k_1) \cosh \theta \sinh \theta = 0$ and $\kappa_n = k_1 \cosh^2 \theta - \sinh^2 \theta = 0$. From the third equation of (2.16),

$$\kappa_g'' = 0$$

The first integral is $\kappa_g' = \text{const.}$

The boundary conditions of (2.16), $\kappa_g'(l) = 0$. Thus $\kappa_g = 0$.

Example 3.2. An arc on first kind helicoid for $\Theta = 0$ in G_1^3 , is elastic line.

Proof. On first kind helicoid, $\kappa_n = 0$ and $\tau_g = 0$. Thus (2.16) is satisfied..

References

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